

RESOLVING THE STRONG CP PROBLEM WITH THE EDM PORTFOLIO*

Jordy de Vries

University of Amsterdam & Nikhef

Based on recent work with Patrick Draper, Kaori Fuyuto, Jonathan Kozaczuk, Ben Lillard

And a lot of older work with the FZJ+Bonn crew (Andreas and Ulf and others)



*Title does not reflect contents of this talk

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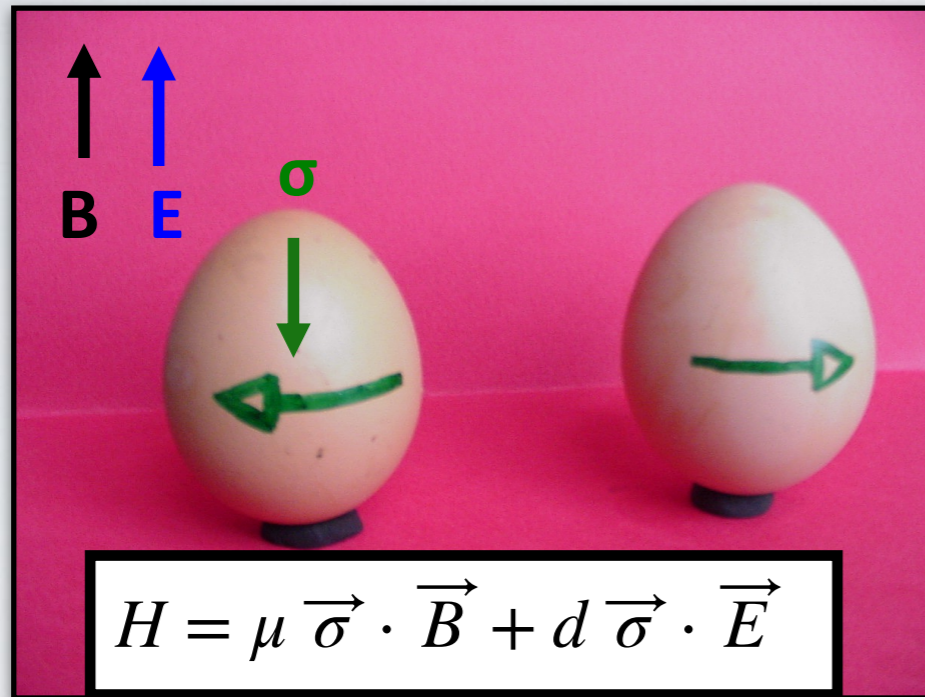


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Electric dipole moments | 0 |



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Diagram showing two eggs on a pink background. The left egg has a green arrow pointing down labeled σ and a green arrow pointing left. The right egg has a green arrow pointing right. Above the eggs, a black arrow labeled B points up, and a blue arrow labeled E points up.

$H = \mu \vec{\sigma} \cdot \vec{B} + d \vec{\sigma} \cdot \vec{E}$

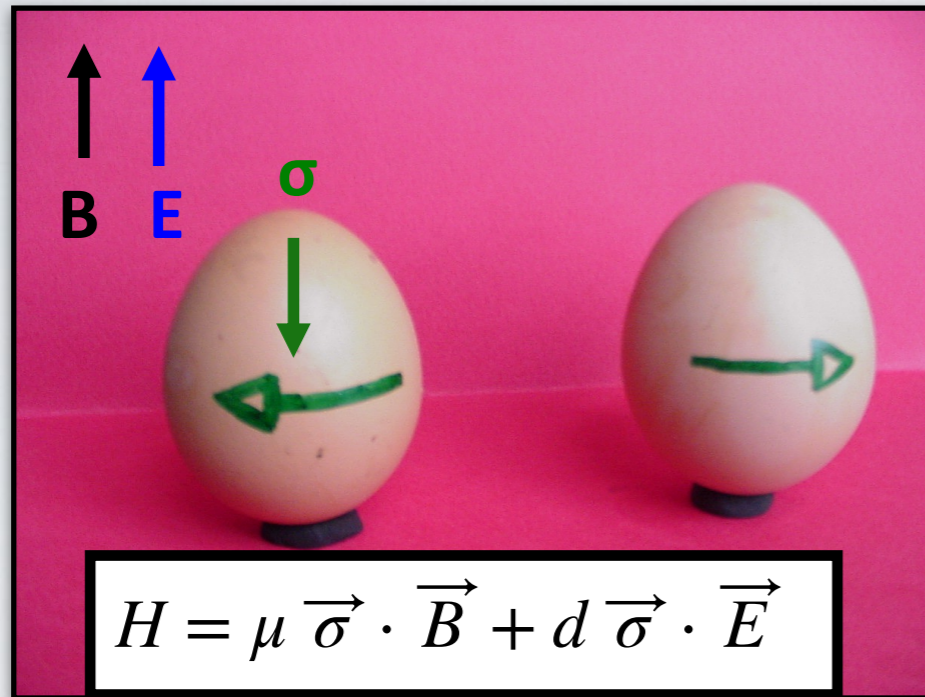
T/CP
transformation



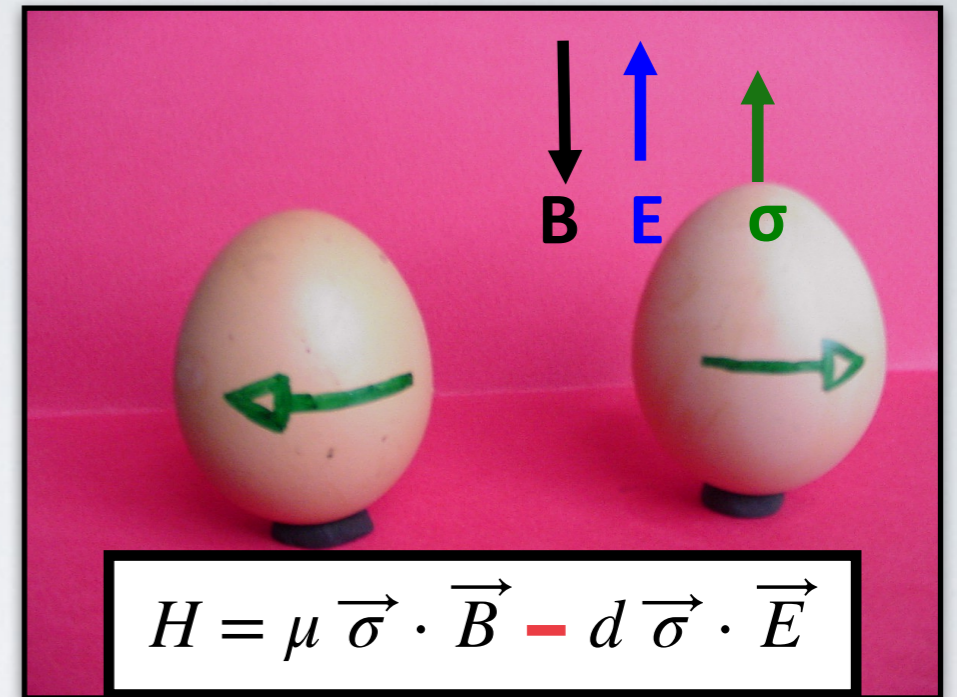
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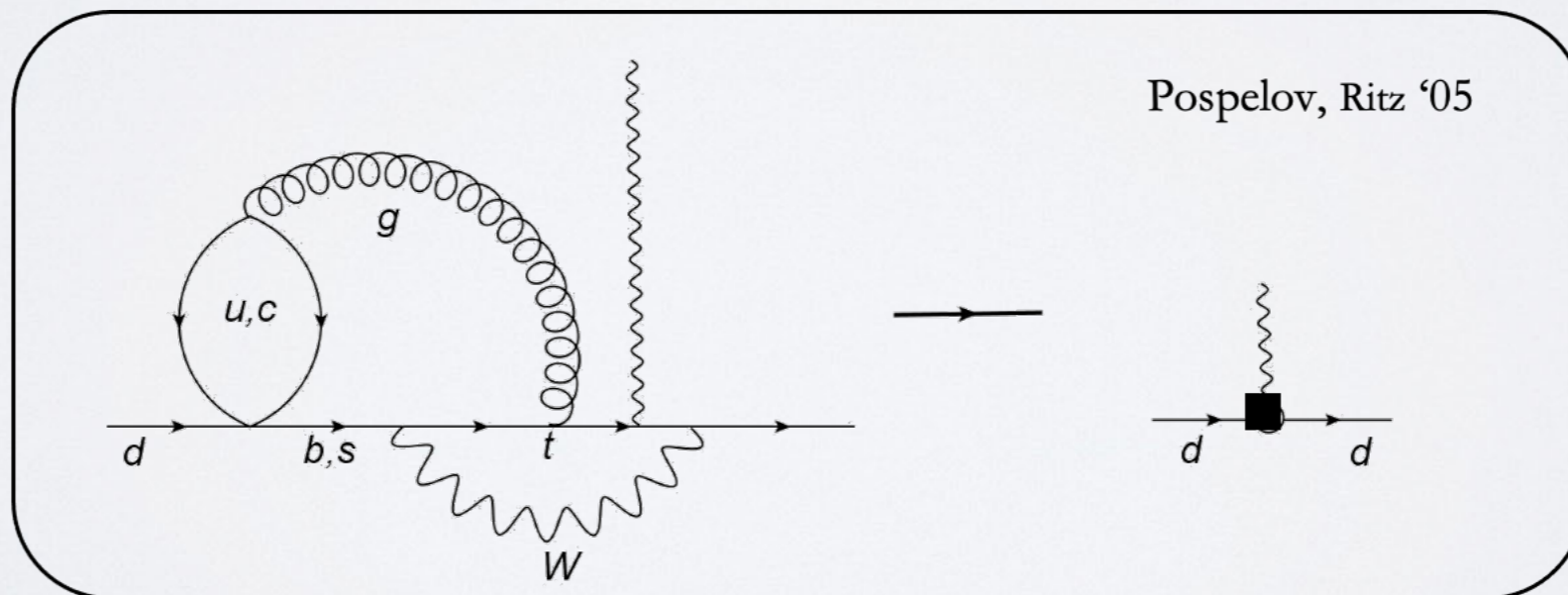
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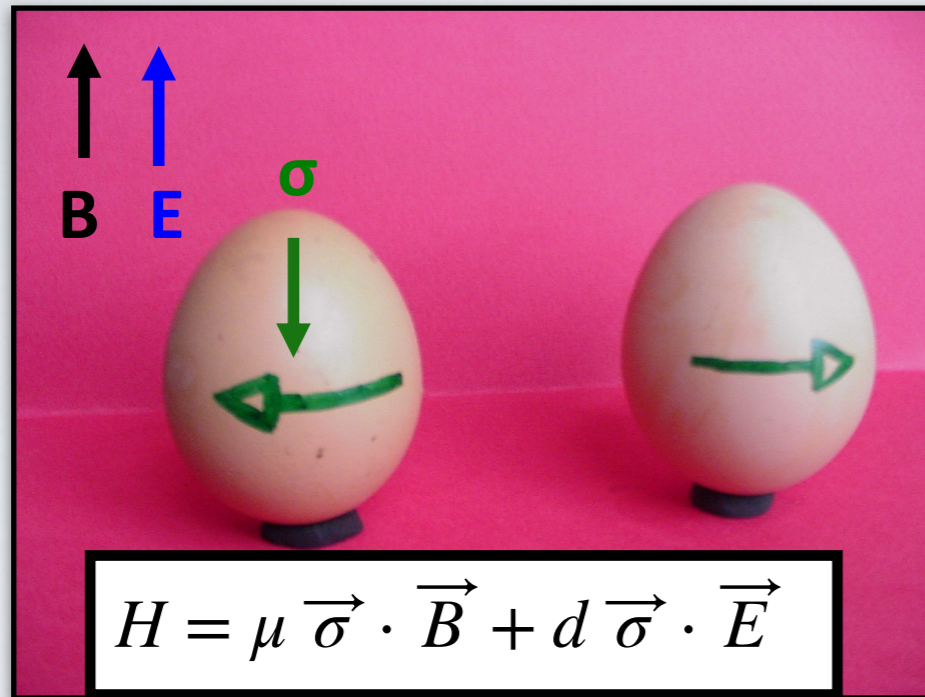
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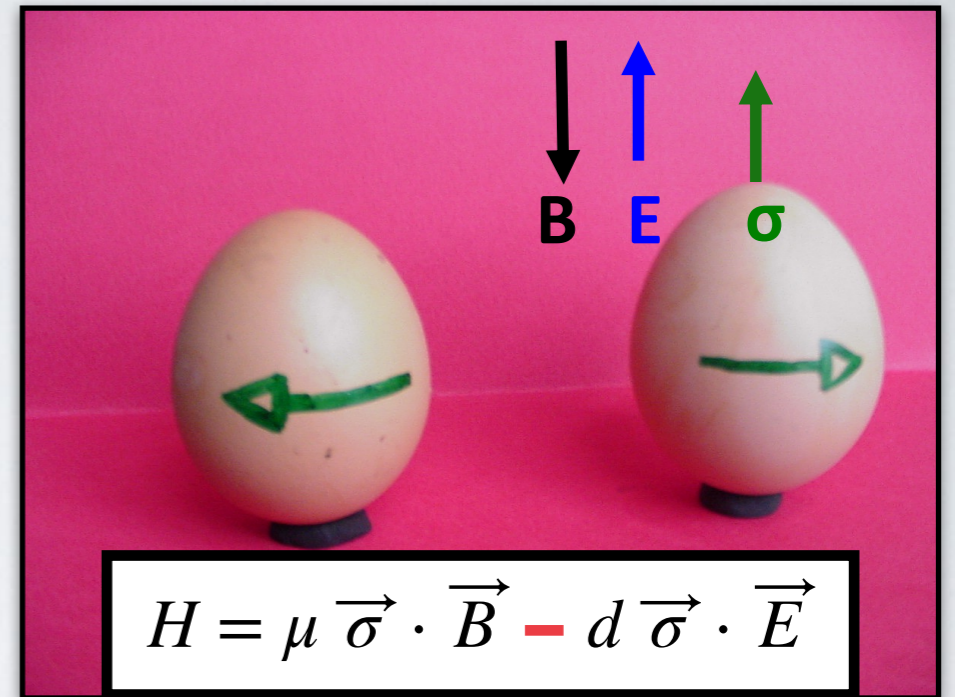
- EDMs from CKM phase only appear at high-loop level and are very suppressed !



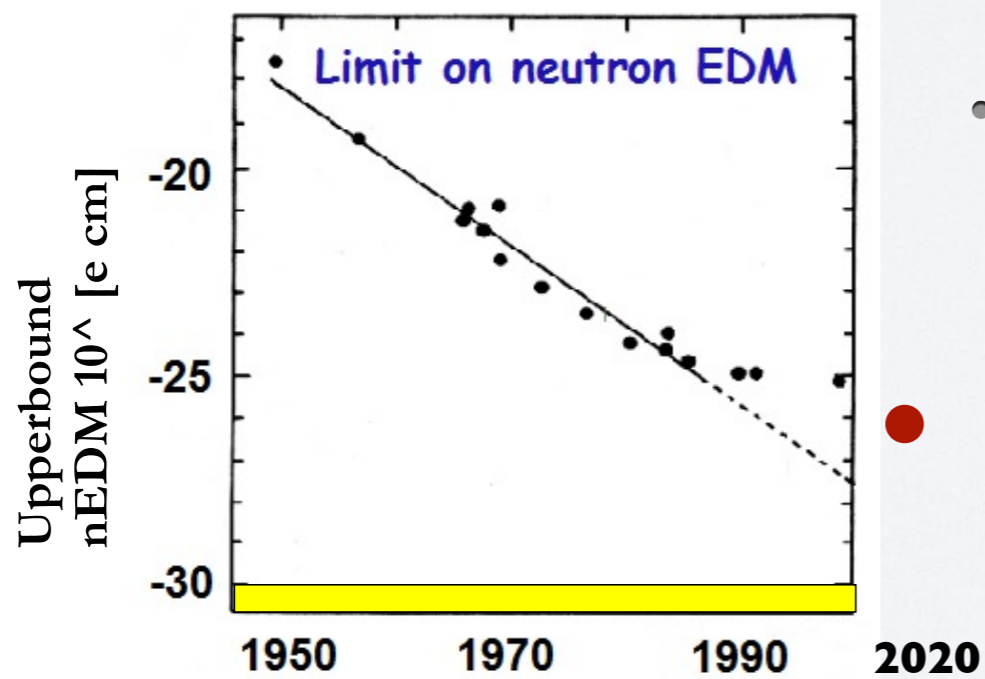
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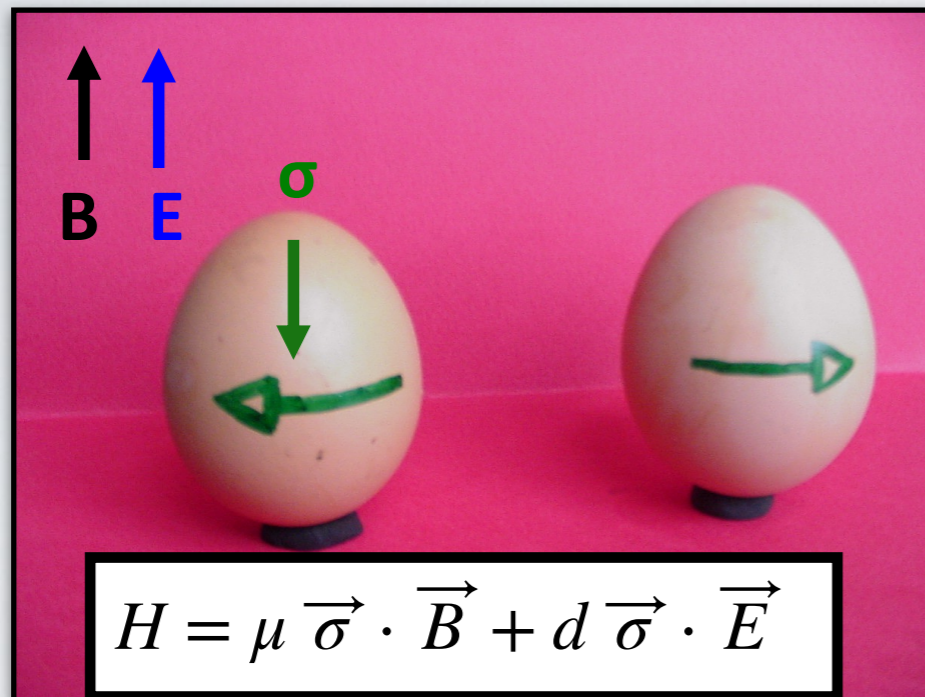


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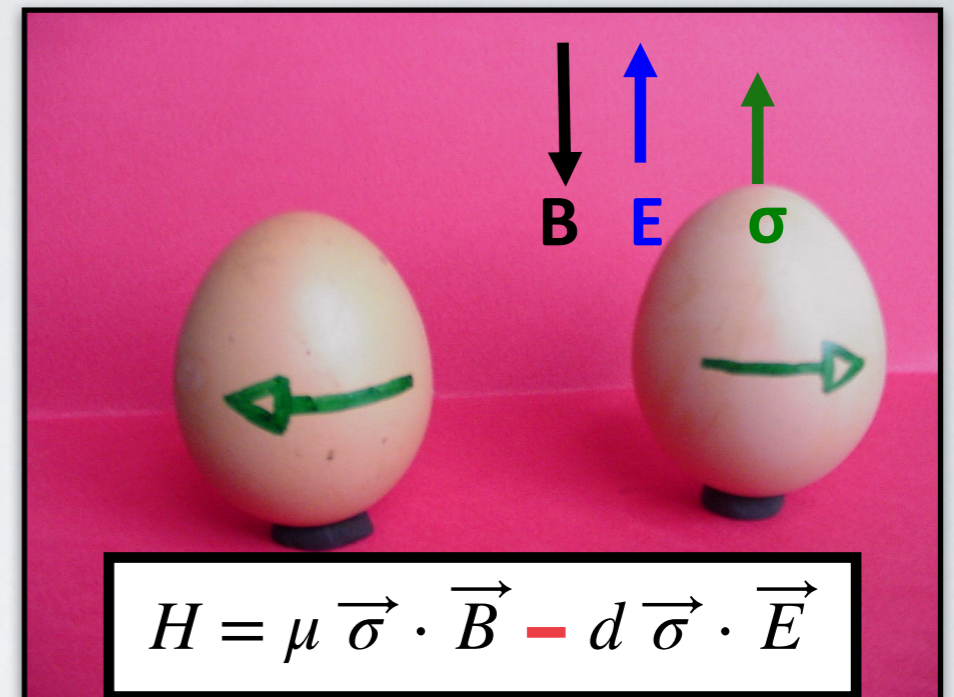


- SM prediction essentially out of reach

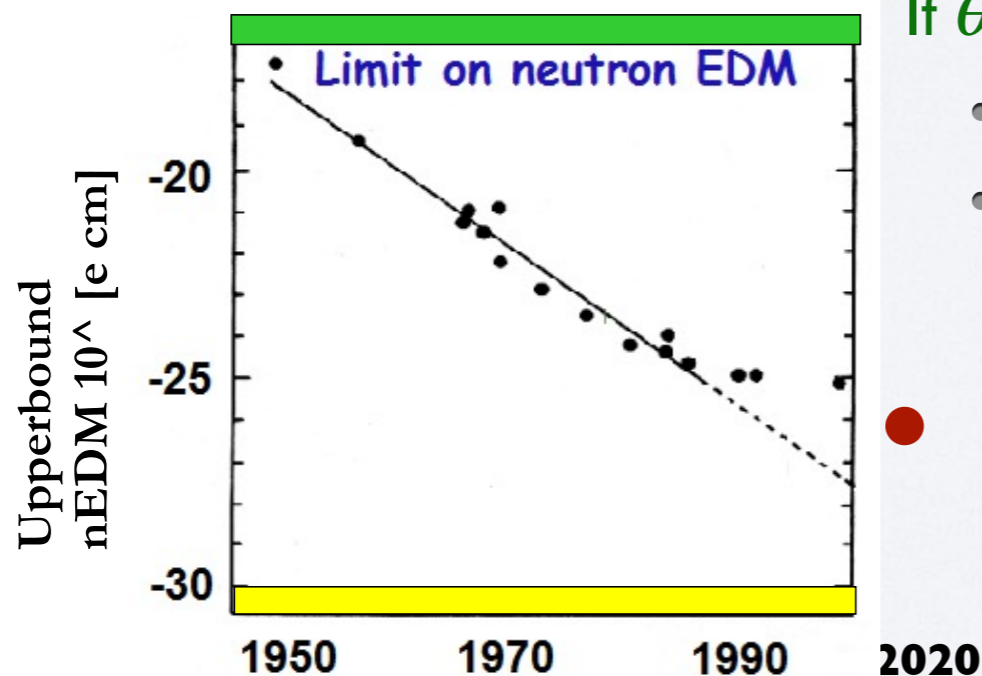
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T/CP transformation



- EDMs from CKM phase only appear at high-loop level and are very suppressed !



If $\theta \sim 1$

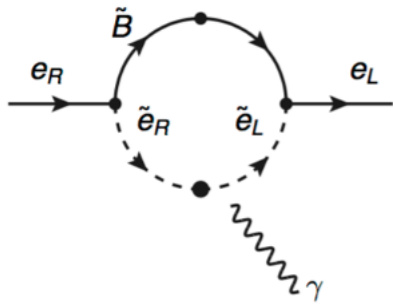
- SM prediction essentially out of reach
- EDMs can still arise from the QCD theta term

$$\mathcal{L}_\theta \sim \bar{\theta} \epsilon^{\mu\nu\alpha\beta} G_{\mu\nu}^a G_{\alpha\beta}^a$$

- Strong CP problem: $\theta < 0.0000000001$
- Sparked a lot of debate and theorizing

Electric dipole moments | 0 |

Example 1:
Bino-Higgsino loop contribution
to the electron EDM



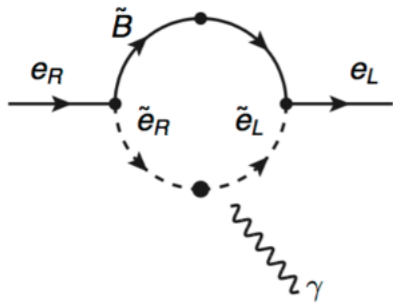
- Many BSM models: EDMs at zero-, one-, or two-loop

$$d_f \left(\frac{\alpha_{em}}{\pi} \right)^n \frac{m_e}{\Lambda^2} \sin \phi_{CPV}$$

- If phase $\sim O(1)$, then $\Lambda > 30 \text{ TeV}$ ($n=1$), or $\Lambda > 2 \text{ TeV}$ ($n=2$)

Electric dipole moments | 0 |

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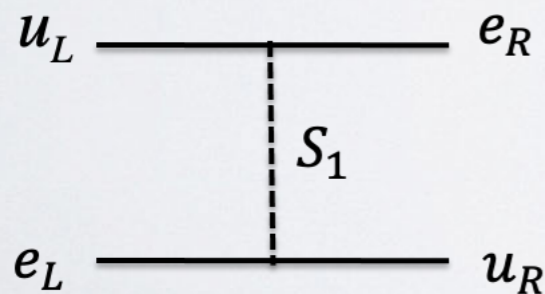
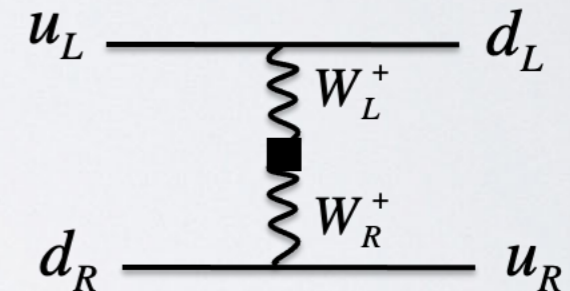


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- If phase $\sim O(1)$, then $\Lambda > 30 \text{ TeV}$ ($n=1$), or $\Lambda > 2 \text{ TeV}$ ($n=2$)

- Certain models EDMs are induced without loop suppression !
- For example, in left-right symmetric models:
- CP-odd four-quark operators induce hadronic EDMs



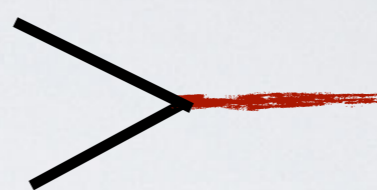
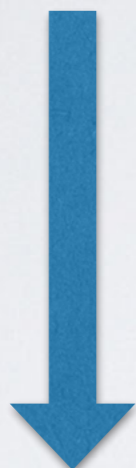
- Leptoquarks can induce CP-odd electron-quark interactions
- Induce atomic/molecular EDMs

- Tree-level CPV leads to $\Lambda > 100\text{-}10000 \text{ TeV}$ if phases are $O(1)$

EDMs are low-energy experiments

Energy

Λ

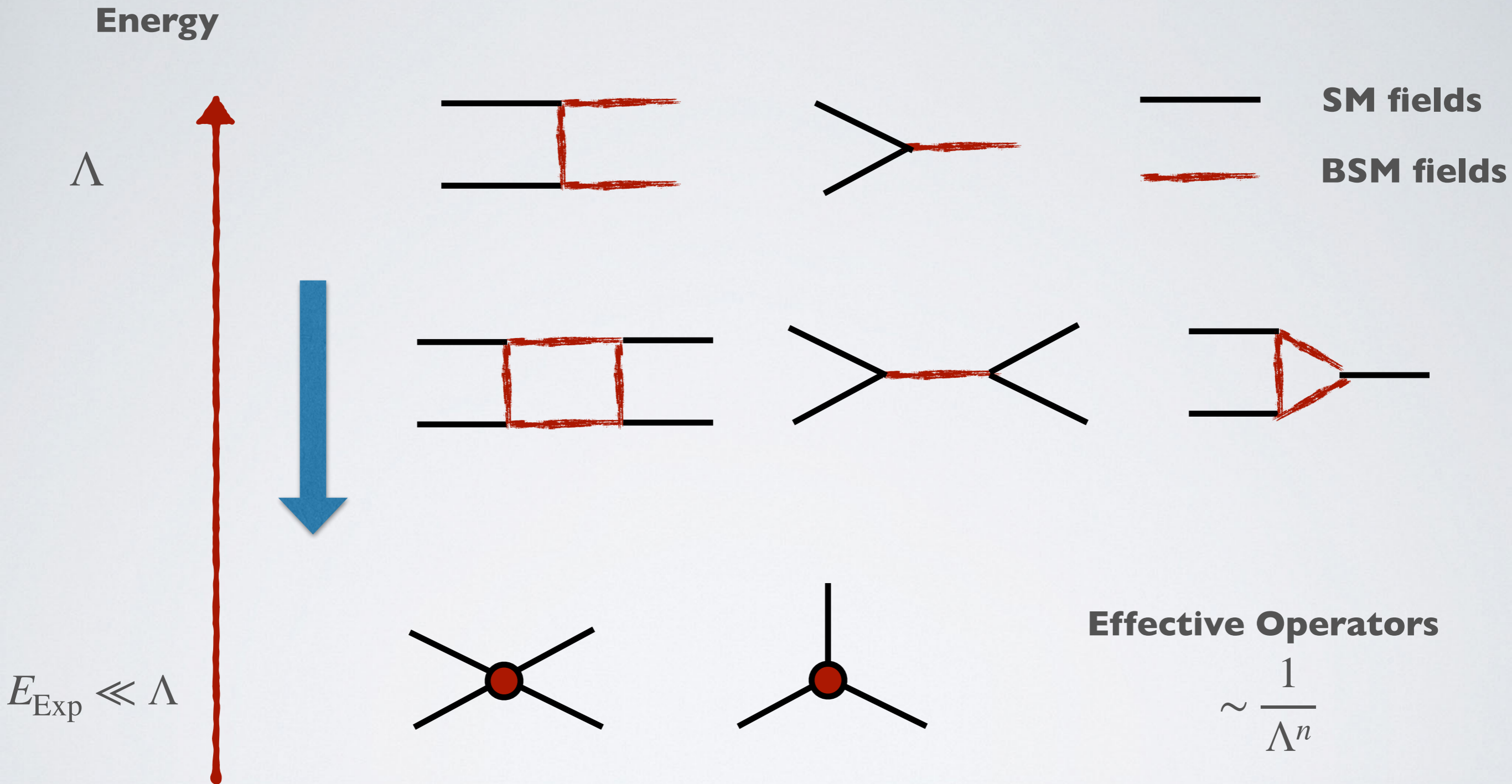


SM fields

BSM fields



EDMs are low-energy experiments



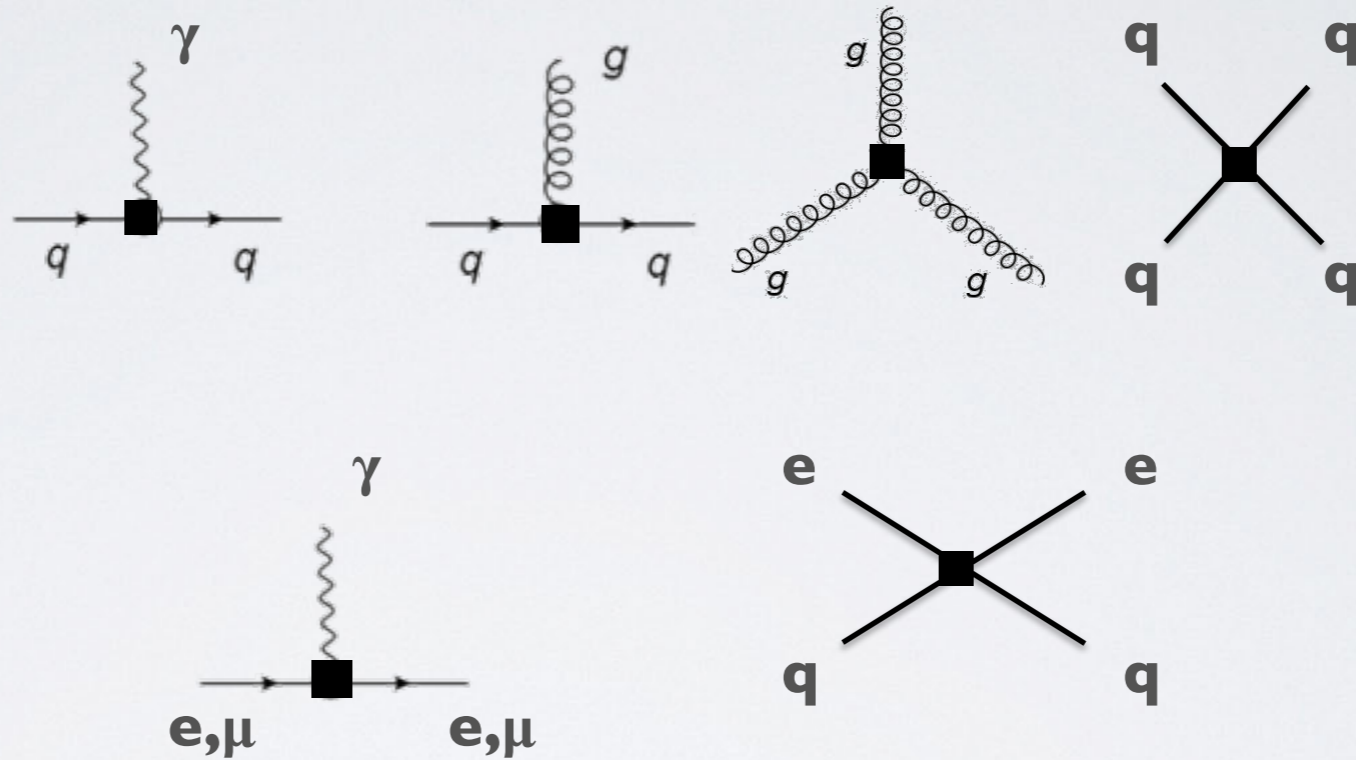
Effects of heavy BSM fields capture by local effective operators

For CP violation relevant operators start at dimension six

Strong CP violation

- Large number of **CP-odd** and **flavor-diagonal** dim-6 operators (unlike Standard Model)
- At energies around a few GeV: handful of operators left

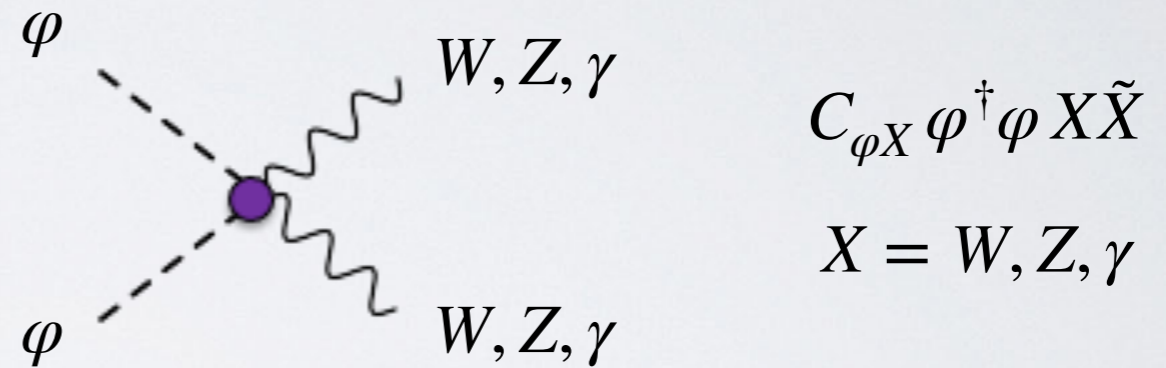
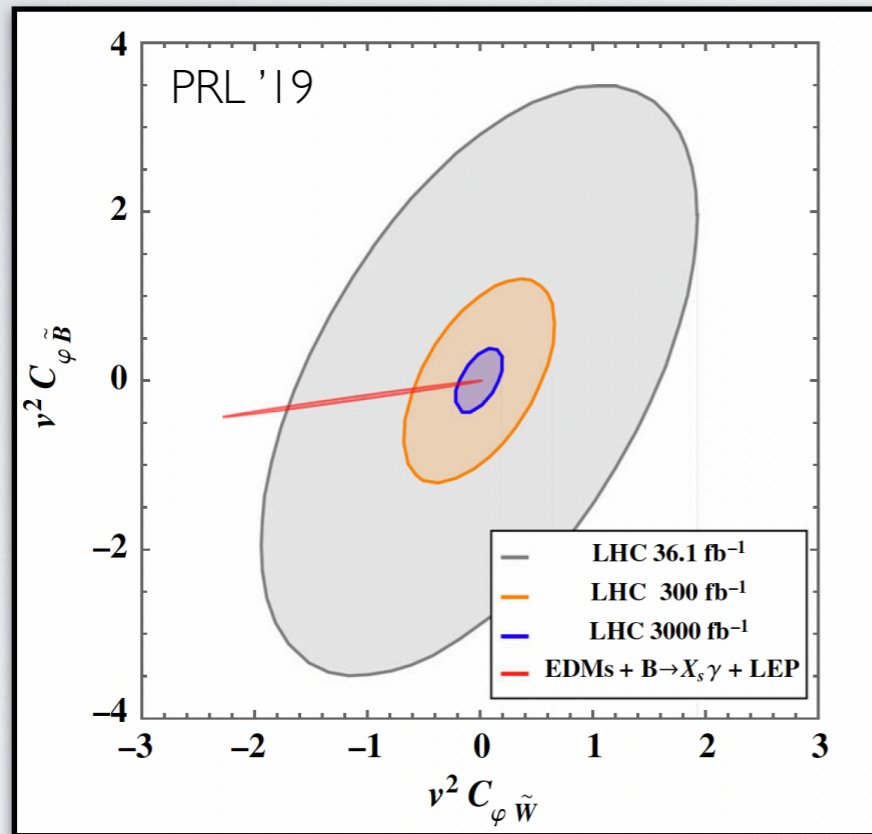
$$\mathcal{L}_\theta \sim \bar{\theta} \epsilon^{\mu\nu\alpha\beta} G_{\mu\nu}^a G_{\alpha\beta}^a$$



- **Induce electric dipole moments of leptons, hadrons, nuclei, atoms, molecules**

CP violation in SM-EFT

- Large number of **CP-odd** and **flavor-diagonal** dim-6 operators (unlike Standard Model)
- At energies around a few GeV: handful of operators left



- **Rich phenomenology of EDMs of leptons, hadrons, nuclei, atoms, molecules**
- Interesting complementarity with collider program. **Example: CP-violation in Higgs sector is best tested by combining LHC + flavor + EDMs.**
- **Direct impact for viability of electroweak baryogenesis**

Experimental searches

	System	Group	Limit	C.L.	Value	Year
e	²⁰⁵ Tl	Berkeley	1.6×10^{-27}	90%	$6.9(7.4) \times 10^{-28}$	2002
	YbF	Imperial	10.5×10^{-28}	90	$-2.4(5.7)(1.5) \times 10^{-28}$	2011
	ThO	ACME	1.1×10^{-29}	90	$4.3(3.1)(2.6) \times 10^{-30}$	2018
	HfF ⁺	Boulder	1.3×10^{-28}	90	$0.9(7.7)(1.7) \times 10^{-29}$	2017
	n	PSI	1.8×10^{-26}	90	$0.0(1.1)(0.2) \times 10^{-26}$	2020
	¹²⁹ Xe	UMich	4.8×10^{-27}	95	$0.26(2.3)(0.7) \times 10^{-27}$	2019
	¹⁹⁹ Hg	UWash	7.4×10^{-30}	95	$-2.2(2.8)(1.5) \times 10^{-30}$	2016
	²²⁵ Ra	Argonne	1.4×10^{-23}	95	$4(6.0)(0.2) \times 10^{-24}$	2016
	muon	E821 BNL g-2	1.8×10^{-19}	95	$0.0(0.2)(0.9) \times 10^{-19}$	2009

+ planned experiments on other systems such as deuteron, Rn, BaF,

A lot of potential for progress !

A Luxury Problem

**Measurement of a
non-zero EDM**



Non-zero θ term

BSM CP violation

Quantifying the strong CP problem

- **Problem:** Calculate EDMs in terms of the theta angle
- First calculation Crewther et al '79, essentially leading-order Chiral perturbation theory.

$$\mathcal{L}_{QCD} = \mathcal{L}_{kin} - \bar{m}\bar{q}q - \epsilon\bar{m}\bar{q}\tau^3q + m_\star\bar{\theta}\bar{q}i\gamma^5q$$

$$m_\star = \frac{m_u m_d}{m_u + m_d}$$

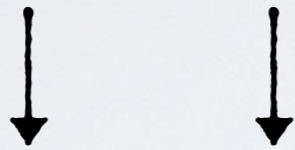
$$\bar{m} = (m_u + m_d)/2$$

$$\epsilon\bar{m} = (m_d - m_u)/2$$

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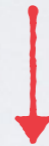
$$\mathcal{L}_{QCD} = \mathcal{L}_{kin} - \bar{m}\bar{q}q - \varepsilon\bar{m}\bar{q}\tau^3q$$



$$\mathcal{L}_{\chi+m} = \mathcal{L}_{\chi} - \frac{m_{\pi}^2}{2}\pi^2 - \delta m_N \bar{N}\tau^3 N$$

**Nucleon mass splitting
(strong part, no EM)**

$$+m_{\star}\bar{\theta}\bar{q}i\gamma^5q$$



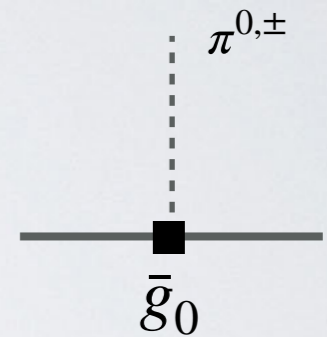
$$+\bar{g}_0\bar{N}\tau \cdot \pi N$$

CP-odd pion-nucleon

$$m_{\star} = \frac{m_u m_d}{m_u + m_d}$$

$$\bar{m} = (m_u + m_d)/2$$

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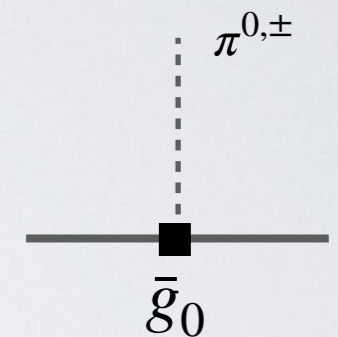
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$$\mathcal{L}_{QCD} = \mathcal{L}_{kin} - \bar{m} \bar{q} q - \varepsilon \bar{m} \bar{q} \tau^3 q \quad \longleftrightarrow \quad + m_\star \bar{\theta} \bar{q} i \gamma^5 q$$

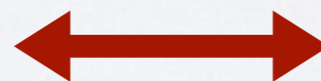
SU_A(2) rotation

$$\mathcal{L}_{\chi+m} = \mathcal{L}_\chi - \frac{m_\pi^2}{2} \pi^2 - \delta m_N \bar{N} \tau^3 N$$

$$+ \bar{g}_0 \bar{N} \tau \cdot \pi N$$



**Nucleon mass splitting
(strong part, no EM)**



CP-odd pion-nucleon

$$\bar{g}_0 = -\frac{\delta m_N}{2f_\pi} \frac{1 - \varepsilon^2}{2\varepsilon} \bar{\theta} = (15.5 \pm 2.5) \cdot 10^{-3} \bar{\theta}$$

δm_N from lattice-QCD

e.g. Borsanyi et al '14

Relation valid up to N²LO corrections

Quantifying the strong CP problem

- **Problem:** Calculate EDMs in terms of the theta angle
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$$d_n = \bar{d}_n(\mu = m_N) - \frac{eg_A\bar{g}_0}{4\pi^2 F_\pi} \left(\ln \frac{m_\pi^2}{m_N^2} - \frac{\pi m_\pi}{2m_N} \right)$$

- The loop part gives $d_n \simeq -2.5 \cdot 10^{-16} e \text{ cm} \longrightarrow \bar{\theta} < 10^{-10}$

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- The loop part gives $d_n \simeq -2.5 \cdot 10^{-16} e \text{ cm}$ $\longrightarrow \bar{\theta} < 10^{-10}$
- Lattice QCD is needed for a full calculation. But no consensus yet it seems.

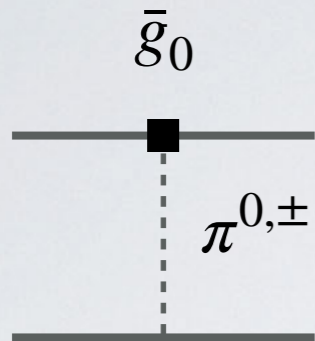
$$d_n = -(1.5 \pm 0.8) \cdot 10^{-16} e \text{ cm} \quad \text{from Shindler et al '19}$$

$$d_n = -(3.9 \pm 1.1) \cdot 10^{-16} e \text{ cm} \quad \text{from Guo et al '15}$$

Neither confirmed by recent calculations from LANL lattice group '21

Other probes of the theta term

CP-odd nuclear force

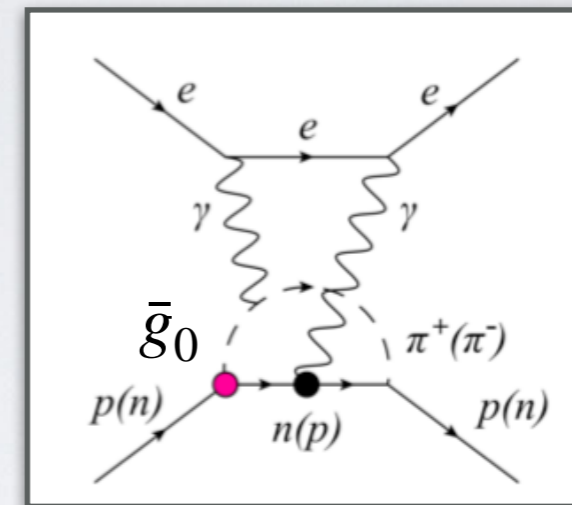


Review: JdV/Meißner '15

Induces EDMs of nuclei and diamagnetic atoms (closed electron shells)

- Diamagnetic atoms (e.g. ^{199}Hg) gives stronger limits but large nuclear uncertainty $\bar{\theta} \sim < 10^{-11}$
- Storage ring experiments would be wonderful (deuteron)
- Polar molecules EDMs not competitive yet, **but experimental progress is rapid.**
Might be the future ! Right now from ThO measurement $\bar{\theta} < 3 \cdot 10^{-8}$

CP-odd nucleon-electron interactions



Flambaum, Pospelov, Ritz, Stadnik '19

Induces EDMs of paramagnetic atoms and molecules

Some musings

Is there really a problem ?

- Not really. **It is just a parameter.** No inconsistencies.
- **Could it have been larger?**
- Seems yes, nothing really changes in the universe if $\theta \sim 0.1$ **No anthropic argument.**

Ubbaldi '08, Inka Hammer '15,

Lee et al '20,

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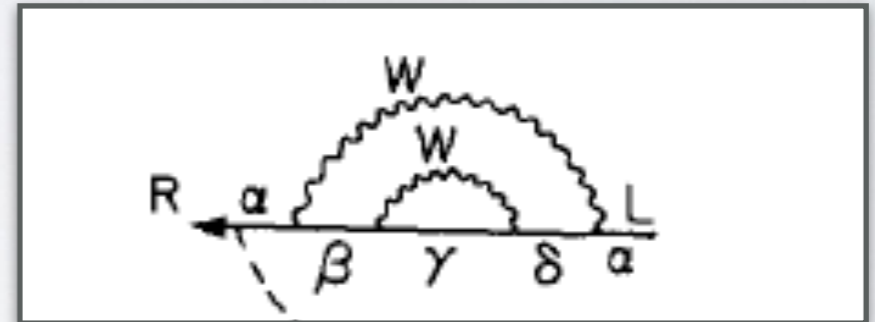
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Is small theta radiatively stable?

- SM has a remarkable property: **theta is technically natural**
- **Ellis/Gaillard '79: tiny CKM contributions**

$$\Delta\bar{\theta} \sim 10^{-16}$$

- This property is lost in generic BSM extensions !



Some musings

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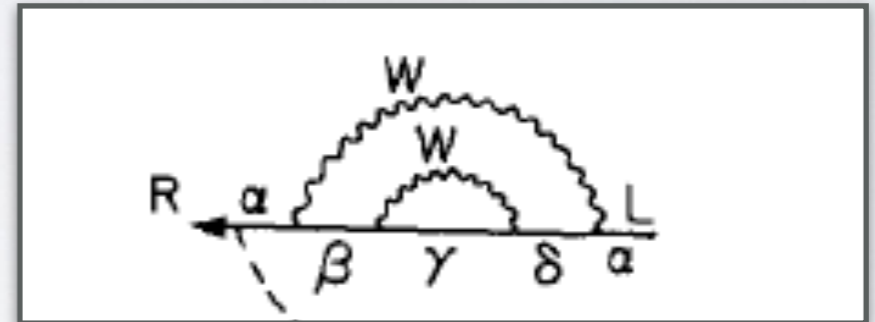
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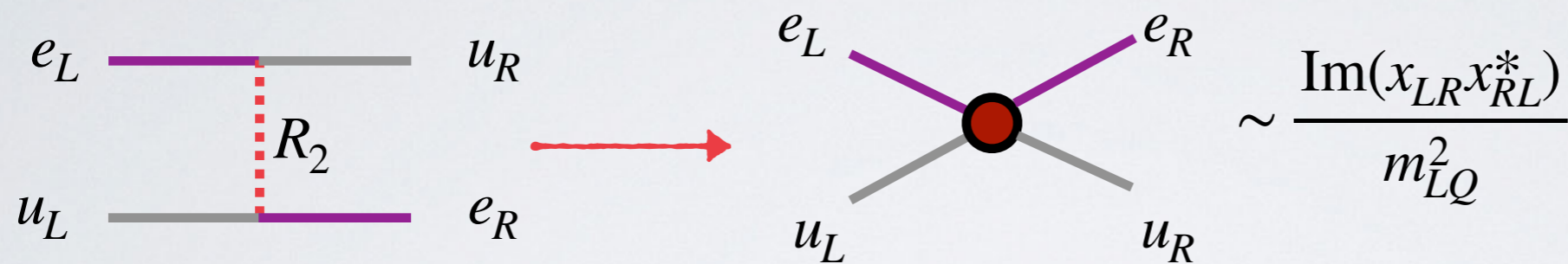
If we do think it is a problem, can we solve it ?

- **UV solutions:** P or CP is a symmetry of UV theory. Break at some scale to generate CKM phase —> Avoid generating a large theta term is not easy!
- **IR solution:** Use a Peccei-Quinn mechanism to dynamically set theta to zero. **AXIONS**
- **Ruled out solution:** massless up quark

The strong CP problem in BSM models

- **In general extension: theta is no longer protected**
- Simplest example: a scalar leptoquark model (so called S_1 LQ for the experts)

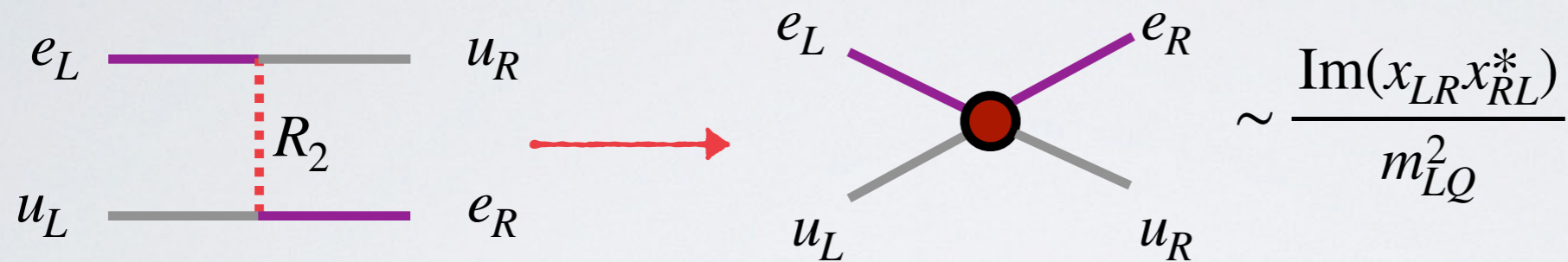
$$\mathcal{L} = R_2 (x_{RL} \bar{u}_R e_L + x_{LR} \bar{u}_L e_R) + \text{h.c.}$$



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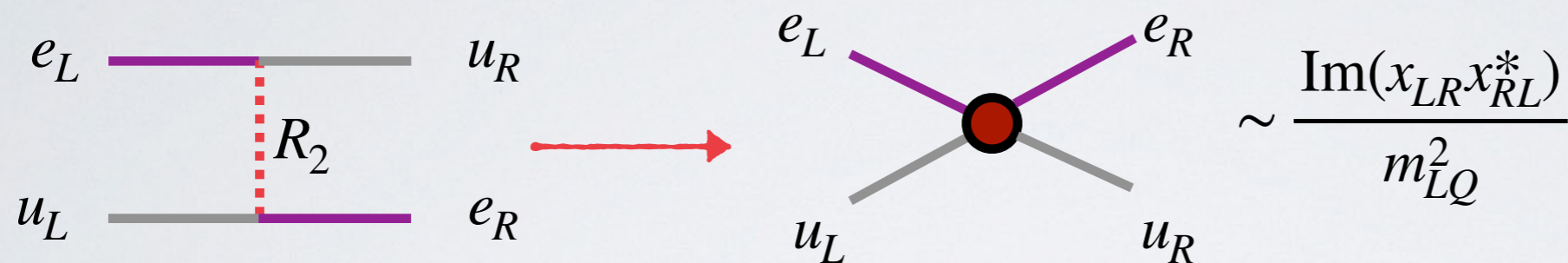
- But also a one-loop contribution to the theta term ! **No decoupling !**

$$\Delta\theta \sim \frac{1}{8\pi^2} \text{Im}[y_u^{-1} x_{RL}^* y_e^\dagger x_{LR}] \sim 10^{-3} \text{Im}[x_{LR} x_{RL}^*]$$

The strong CP problem in BSM models

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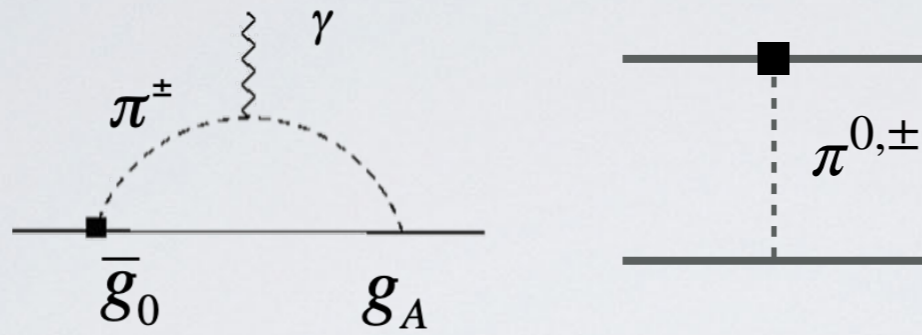


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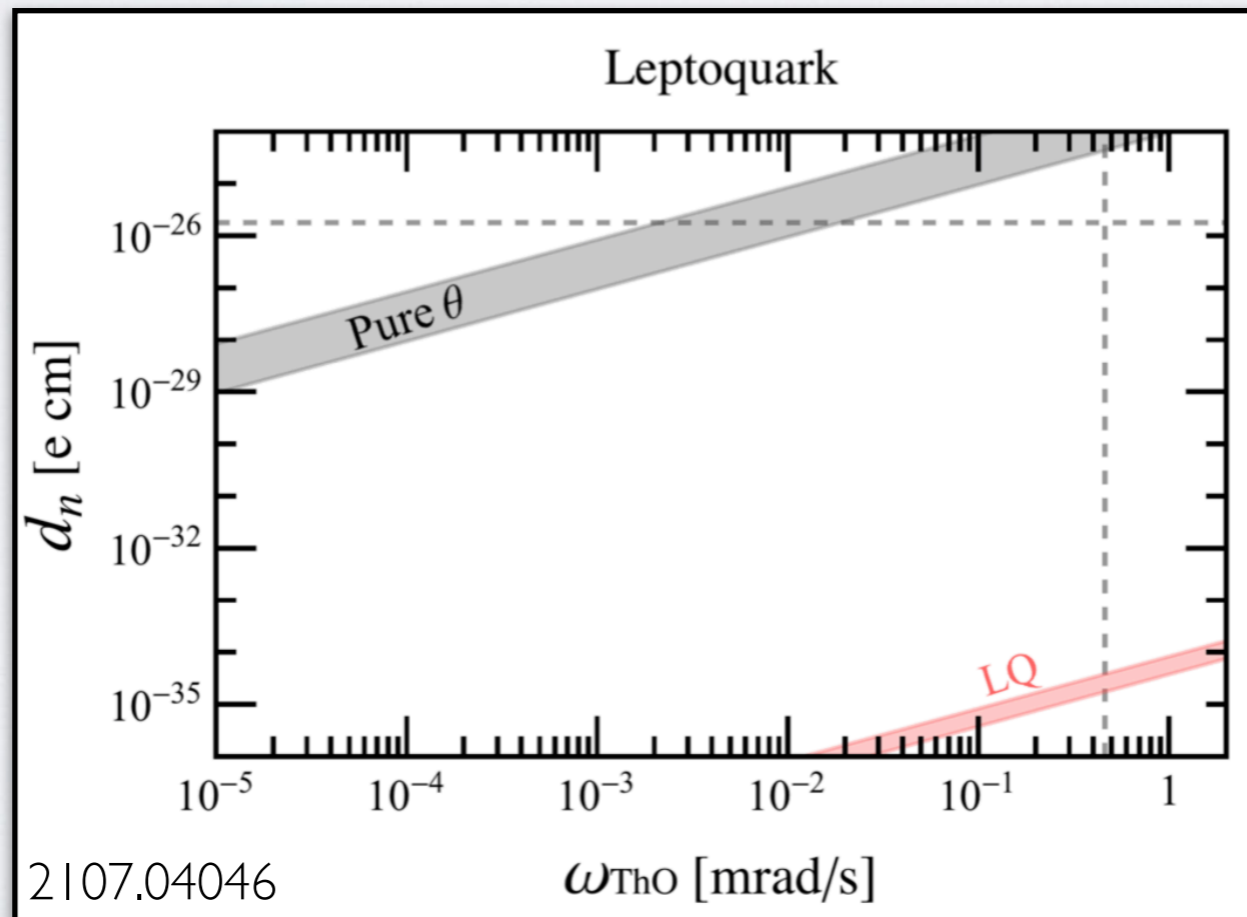
$$\Delta\theta \sim \frac{1}{8\pi^2} \text{Im}[y_u^{-1} x_{RL}^* y_e^\dagger x_{LR}] \sim 10^{-3} \text{Im}[x_{LR} x_{RL}^*]$$

- UV 'solution': Set phase to $< 10^{-7}$. Note that the dim-6 piece vanishes even quicker !
- **EDMs are now dominated by the 'remainder' of theta term.**
- Or Peccei-Quinn mechanism at low energies to effectively set $\theta_0 + \Delta\theta \rightarrow 0$
- But under a PQ: **Dimension-six term sticks around.**
- **EDMs dominated by dimension-six electron-quark interactions.**

- **UV solution:** Low-energy CPV dominated by theta. Neutron + Hg \gg Paramagnetic EDMs



- **PQ mechanism:** Low-energy CPV dominated by dim-six electron-nucleon operators



Can we generalize this?

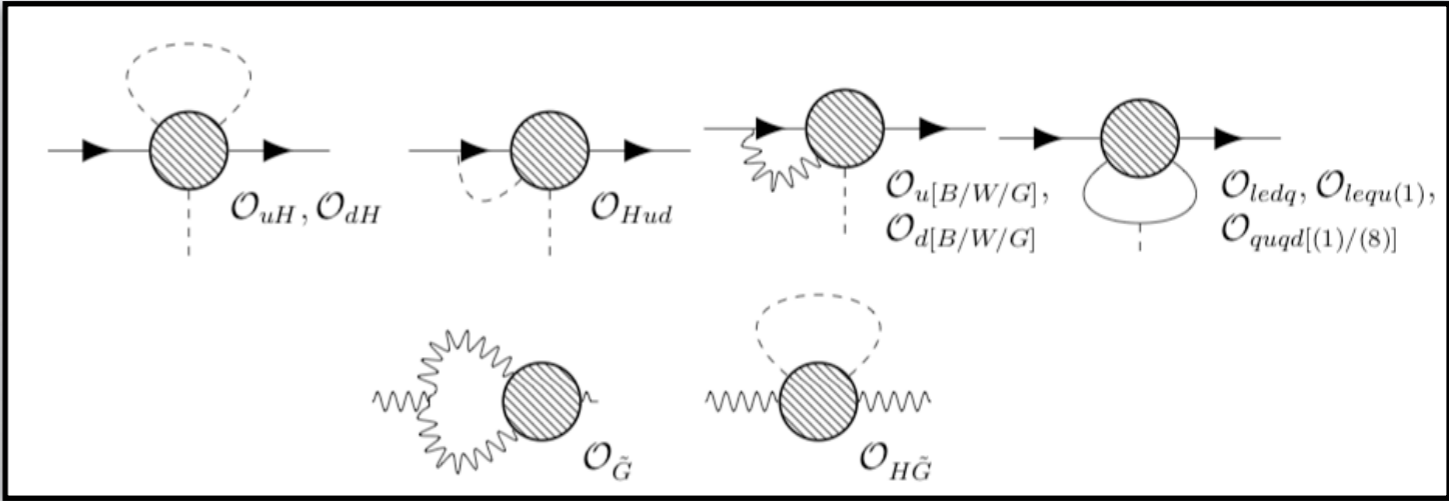
- This was just some stupid model. Is there a more thorough argument?
- Use divergence structure of dimension-six SM-EFT operators.

$$\delta\bar{\theta} = \delta\theta + \delta(\arg \det Y_u) + \delta(\arg \det Y_d)$$

- We study the mixing of EFT operators with theta term

Patrick Draper et al '18

\mathcal{O}_{uH}	$H^\dagger H \overline{Q_{Li}} \tilde{H} u_{Rj}$	\mathcal{O}_{dH}	$H^\dagger H \overline{Q_{Li}} H d_{Rj}$
\mathcal{O}_{dG}	$\overline{Q_{Li}} \sigma^{\mu\nu} T^a d_{Rj} H G_{\mu\nu}^a$	\mathcal{O}_{dW}	$\overline{Q_{Li}} \sigma^{\mu\nu} d_{Rj} \tau^a H W_{\mu\nu}^a$
\mathcal{O}_{dB}	$\overline{Q_{Li}} \sigma^{\mu\nu} d_{Rj} H B_{\mu\nu}$	\mathcal{O}_{uG}	$\overline{Q_{Li}} \sigma^{\mu\nu} T^a u_{Rj} \tilde{H} G_{\mu\nu}^a$
\mathcal{O}_{uW}	$\overline{Q_{Li}} \sigma^{\mu\nu} u_{Rj} \tau^a \tilde{H} W_{\mu\nu}^a$	\mathcal{O}_{uB}	$\overline{Q_{Li}} \sigma^{\mu\nu} u_{Rj} \tilde{H} B_{\mu\nu}$
\mathcal{O}_{Hud}	$i \tilde{H}^\dagger D_\mu H \overline{u_{Ri}} \gamma^\mu d_{Rj}$	$\mathcal{O}_{quqd(1)}$	$\epsilon^{ef} \overline{Q_{Li}^e} u_{Rj} \overline{Q_{Lk}^f} d_{Rl}$
$\mathcal{O}_{quqd(8)}$	$\epsilon^{ef} \overline{Q_{Li}^e} T^a u_{Rj} \overline{Q_{Lk}^f} T^a d_{Rl}$	$\mathcal{O}_{lequ(1)}$	$\epsilon^{ef} \overline{L_{Li}^e} e_{Rj} \overline{Q_{Lk}^f} u_{Rl}$
\mathcal{O}_{ledq}	$\overline{L_{Li}} e_{Rj} \overline{d_{Rk}} Q_{Ll}$	\mathcal{O}_{HG}	$H^\dagger H G_{\mu\nu}^a \tilde{G}^{a\mu\nu}$
$\mathcal{O}_{\tilde{G}}$	$f^{abc} G_{\nu}^{a\mu} G_{\rho}^{b\nu} \tilde{G}^{c\rho}_{\mu}$		



Can we generalize this?

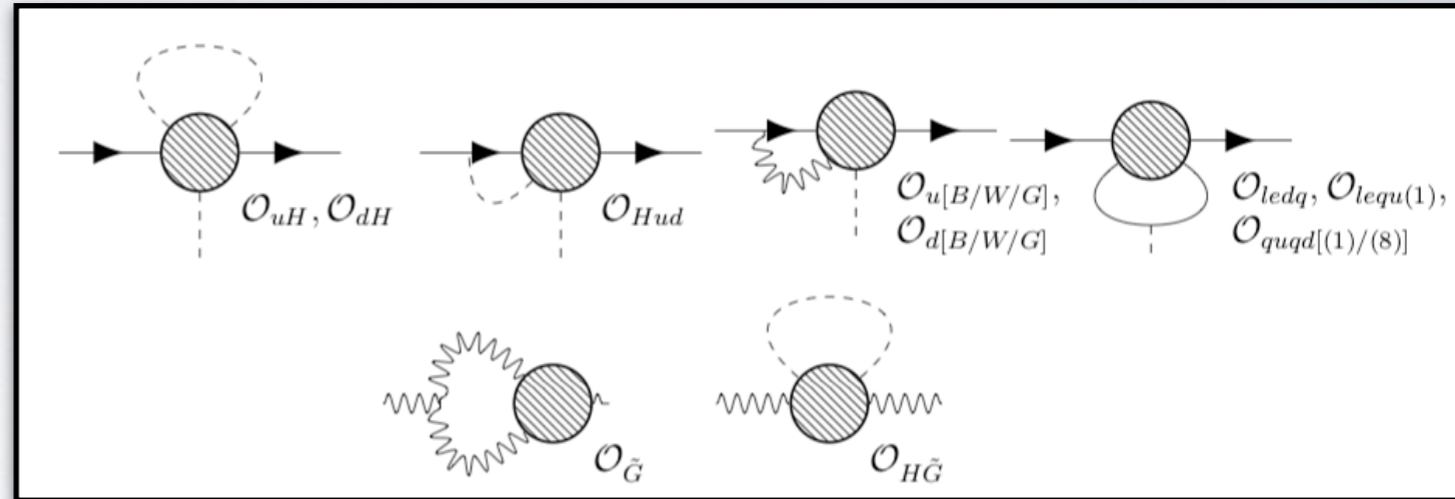
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$$\delta\bar{\theta} = \delta\theta + \delta(\arg \det Y_u) + \delta(\arg \det Y_d)$$

- We study the mixing of EFT operators with theta term

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\mathcal{O}_{uH}	$\frac{H^\dagger H \overline{Q_{Li}} \tilde{H} u_{Rj}}$	\mathcal{O}_{dH}	$\frac{H^\dagger H \overline{Q_{Li}} H d_{Rj}}$
\mathcal{O}_{dG}	$\frac{\overline{Q_{Li}} \sigma^{\mu\nu} T^a d_{Rj} H G_{\mu\nu}^a}{Q_{Li} \sigma^{\mu\nu} d_{Rj} H B_{\mu\nu}}$	\mathcal{O}_{dW}	$\frac{\overline{Q_{Li}} \sigma^{\mu\nu} d_{Rj} \tau^a H W_{\mu\nu}^a}{Q_{Li} \sigma^{\mu\nu} T^a u_{Rj} \tilde{H} G_{\mu\nu}^a}$
\mathcal{O}_{dB}	$\frac{\overline{Q_{Li}} \sigma^{\mu\nu} d_{Rj} H B_{\mu\nu}}{Q_{Li} \sigma^{\mu\nu} u_{Rj} \tau^a \tilde{H} W_{\mu\nu}^a}$	\mathcal{O}_{uG}	$\frac{\overline{Q_{Li}} \sigma^{\mu\nu} T^a u_{Rj} \tilde{H} G_{\mu\nu}^a}{Q_{Li} \sigma^{\mu\nu} u_{Rj} \tilde{H} B_{\mu\nu}}$
\mathcal{O}_{uW}	$\frac{\overline{Q_{Li}} \sigma^{\mu\nu} u_{Rj} \tau^a \tilde{H} W_{\mu\nu}^a}{i \tilde{H}^\dagger D_\mu H \overline{u_{Ri}} \gamma^\mu d_{Rj}}$	\mathcal{O}_{uB}	$\frac{\overline{Q_{Li}} \sigma^{\mu\nu} u_{Rj} \tilde{H} B_{\mu\nu}}{\epsilon^{ef} \overline{Q_{Li}^e} T^a u_{Rj} \overline{Q_{Lk}^f} T^a d_{Rl}}$
\mathcal{O}_{Hud}	$\frac{i \tilde{H}^\dagger D_\mu H \overline{u_{Ri}} \gamma^\mu d_{Rj}}{\epsilon^{ef} \overline{Q_{Li}^e} T^a u_{Rj} \overline{Q_{Lk}^f} T^a d_{Rl}}$	$\mathcal{O}_{quqd(1)}$	$\frac{\epsilon^{ef} \overline{Q_{Li}^e} T^a u_{Rj} \overline{Q_{Lk}^f} T^a d_{Rl}}{\overline{L_{Li}} e_{Rj} \overline{d_{Rk}} Q_{Ll}}$
$\mathcal{O}_{quqd(8)}$	$\frac{\epsilon^{ef} \overline{Q_{Li}^e} T^a u_{Rj} \overline{Q_{Lk}^f} T^a d_{Rl}}{f^{abc} G_{\nu}^{a\mu} G_{\rho}^{b\nu} \tilde{G}^{c\rho}_{\mu}}$	$\mathcal{O}_{lequ(1)}$	$\frac{\epsilon^{ef} \overline{L_{Li}^e} e_{Rj} \overline{Q_{Lk}^f} u_{Rl}}{H^\dagger H G_{\mu\nu}^a \tilde{G}^{a\mu\nu}}$
\mathcal{O}_{ledq}	$\frac{\overline{L_{Li}} e_{Rj} \overline{d_{Rk}} Q_{Ll}}{f^{abc} G_{\nu}^{a\mu} G_{\rho}^{b\nu} \tilde{G}^{c\rho}_{\mu}}$	\mathcal{O}_{HG}	$\frac{H^\dagger H G_{\mu\nu}^a \tilde{G}^{a\mu\nu}}{f^{abc} G_{\nu}^{a\mu} G_{\rho}^{b\nu} \tilde{G}^{c\rho}_{\mu}}$
$\mathcal{O}_{\tilde{G}}$			



- Quadratic divergences (i.e. cut-off scheme) signal unsuppressed threshold corrections to theta
- Of course: cannot calculate threshold corrections within the EFT ! **Just a diagnostic tool**

$$16\pi^2 \delta\bar{\theta} \sim 16\pi^2 \left(\frac{2}{g_s^2} c_{HG} - \frac{9}{2g_s} c_{\tilde{G}} \right) + \text{Im Tr}[Y_d^{-1} (3c_{dH} + g'c_{dB} - 18gc_{dW} - 16g_sc_{dG})] + \text{Im Tr}[Y_u^{-1} (3c_{uH} - 5g'c_{uB} - 18gc_{uW} - 16g_sc_{uG})] + \text{Im Tr}[(Y_d^{-1} Y_u + Y_d^\dagger (Y_u^\dagger)^{-1}) c_{Hud}] + \text{Im}[2c_{lequ(1)}^{mnij} Y_e^{\dagger nm} (Y_u^{-1})^{ji} - 2c_{ledq}^{*mnij} Y_e^{mn} (Y_d^{-1})^{ij}] + \text{Im}[(6c_{quqd(1)}^{mnij} + c_{quqd(1)}^{inmj} + \frac{4}{3}c_{quqd(8)}^{inmj})(Y_u^{\dagger nm} (Y_d^{-1})^{ji} + Y_d^{\dagger ji} (Y_u^{-1})^{nm})].$$

Can we generalize this?

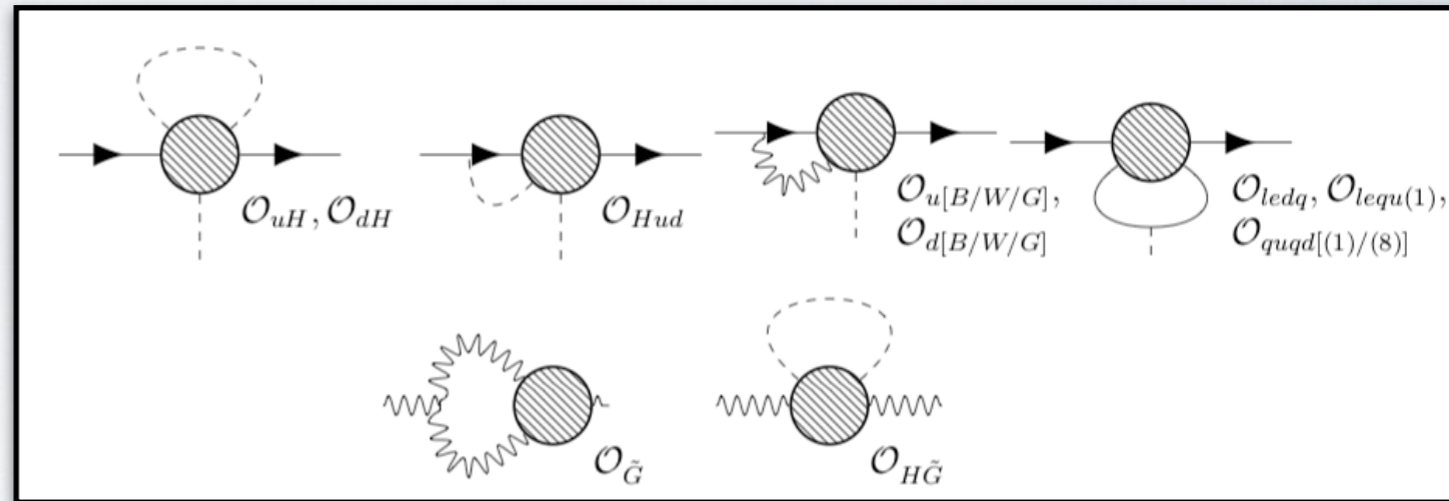
- This was just some stupid model. Is there a more thorough argument?
- Use divergence structure of dimension-six SM-EFT operators.

$$\delta\bar{\theta} = \delta\theta + \delta(\arg \det Y_u) + \delta(\arg \det Y_d)$$

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- Most dim-six operators relevant for EDMs, mix 'quadratically' with theta
- **If dim-6 dominate EDMs at low energy, then where is the dim-4 theta term?**
- **Should be resolved in the IR: a Peccei-Quin mechanism.**
- Deviations from the theta-term EDM pattern provide a hint for an IR solution (PQ mechanism)

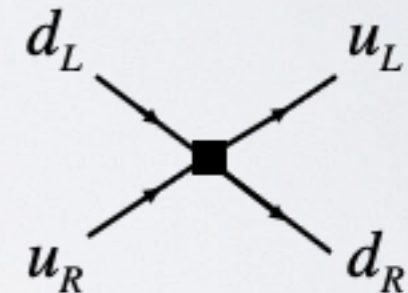
Testing the EFT arguments

JdV et al '21

- Studied a broad class of models (SUSY, 2HDM, leptoquarks, left-right symmetry)
- In all cases, when a dimension-six operator appears: also large threshold correction to theta
 1. Suppress theta threshold at UV scale: all CPV gets suppressed and **theta dominates**
 2. Eliminate theta in IR by Peccei-Quinn: **Dimension-six CPV dominates.**

- Example: minimal left-right symmetric model.

$$\Delta\bar{\theta} \sim \sin\alpha \quad \mathcal{L}_{6,CPV} \sim \frac{\sin\alpha}{m_{W_R}^2} \left[(\varphi^\dagger D_\mu \varphi) \bar{u}_R \gamma^\mu d_R - \text{h.c.} \right]$$



- Dim-6 operator leads to CP-odd isospin-breaking four-quark operator

- **This has a very different EDM pattern** than theta!

- Can only be relevant if somehow $\Delta\bar{\theta}$ gets removed but not setting $\sin\alpha \rightarrow 0$

Testing the EFT arguments

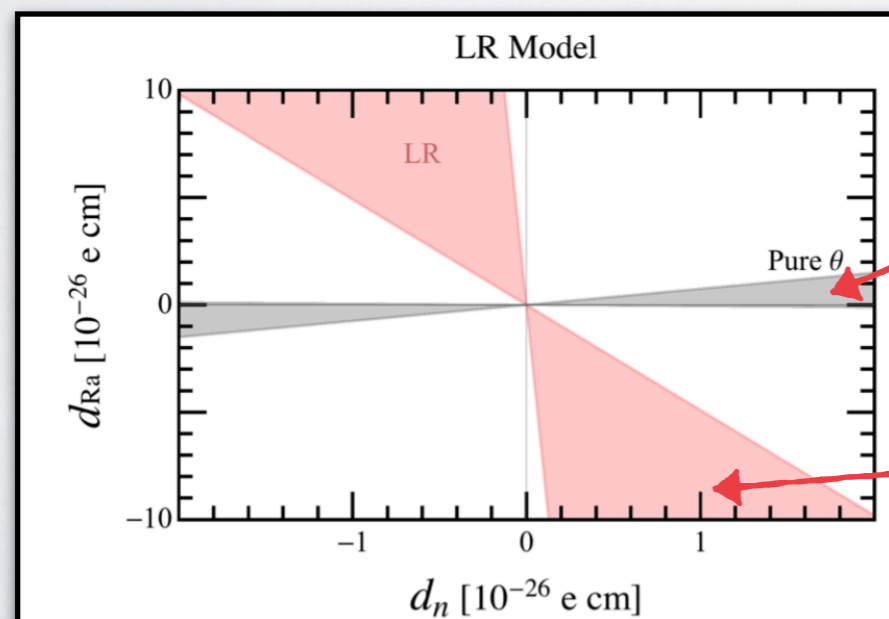
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Left-right model with UV solution

Left-right model with Peccei-Quinn

Testing the EFT arguments

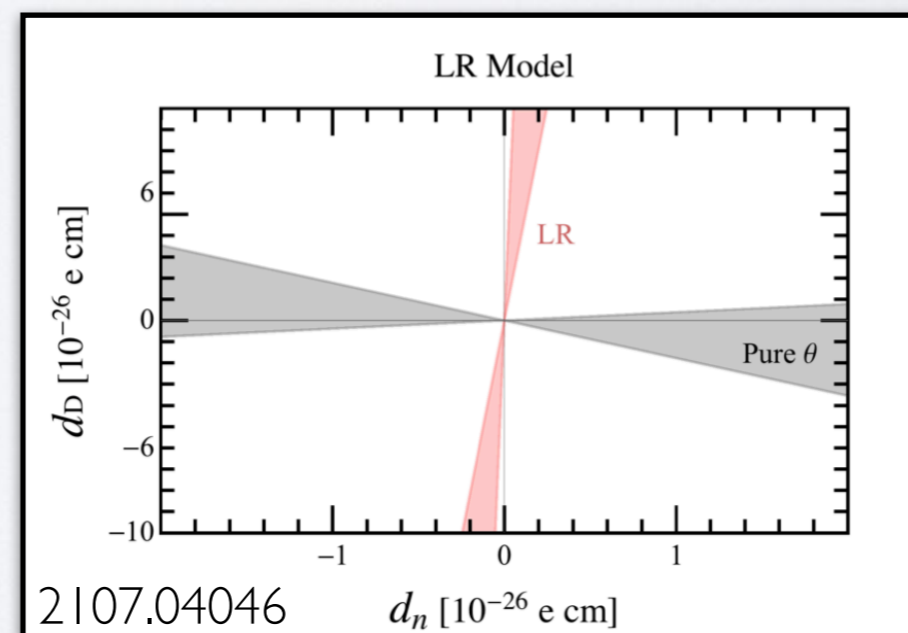
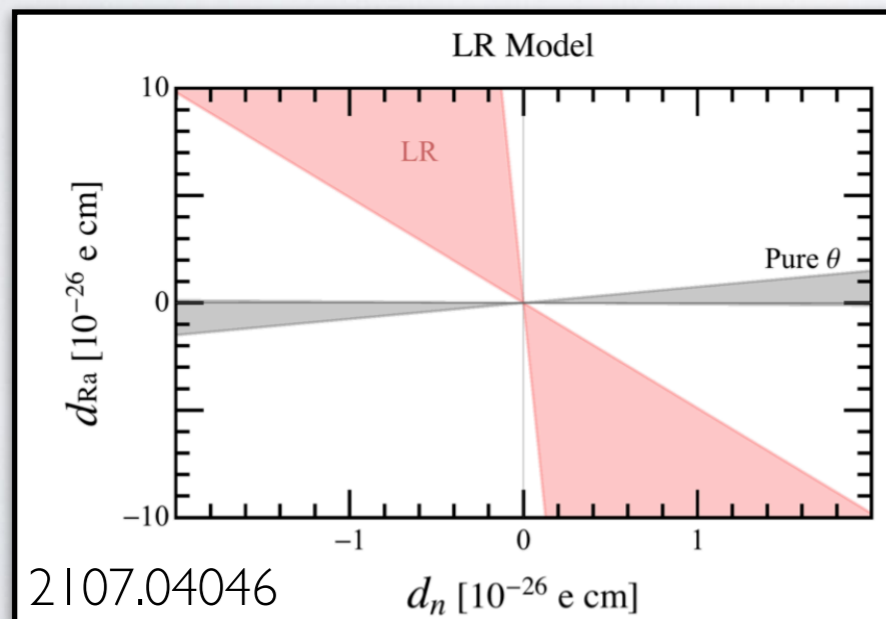
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Conclusions/Summary

- EDMs are powerful ways to look for new CP violation
- Sensitive to dimension-six sources up to thousands of TeV (depending on operator)
- **Last decade, a lot of theory improvements to calculate EDMs (EFT, lattice)**
- Problem: there exists a dim-4 CPV term in SM (theta term)
- **Technically natural in SM, but not in generic BSM models**
- Reflected by mixing pattern of SM-EFT operators with theta
- Models that generate dim-6 operators also induce an unsuppressed theta
- Conclusion: pattern of EDMs that are inconsistent with theta term imply an IR mechanism of the strong CP problem → **Peccei-Quinn mechanism is only game in town.**
- So EDM measurements could hint towards a PQ mechanism without **seeing actual axions.**
- But naturalness arguments are dangerous.... So this would just be a hint, not a proof.

Work in progress/Outlook

- The combination of dimension-six CPV and Peccei-Quinn leads to CPV axion interactions.
- In addition, to the operators discussed by Thomas yesterday
- CPV axion couplings take the simple form

$$\mathcal{L}_{a,CPV} = g_s^e a \bar{e}e + g_s^N a \bar{N}N + g_s^\pi a \pi^2$$

- In progress: calculate these couplings in SMEFT + ChPT
- Compare EDM limits on SMEFT couplings to limits on axion couplings (fifth-force experiments)

